

# Remarks on, CFTs, Radial Lattice Quantization\* and Graphene\*\*!

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Boston University, May 10, 2014

**PPCM:**

**Field Theoretic Computer Simulations  
for Particle Physics and Condensed Matter**

\*RCB, G. Fleming, H. Neuberger & M. Cheng

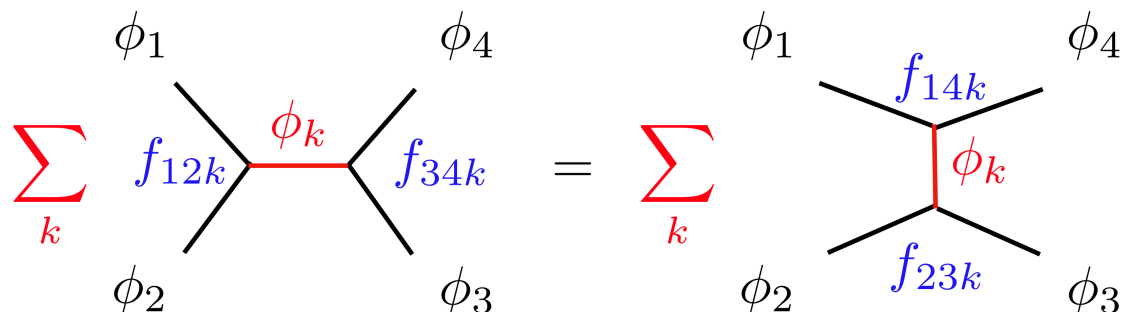
\*\*RCB, C. Rebbi, D. Schaich & M. Cheng

# CFT are highly constrained

$$\langle \phi(x_1)\phi(x_2) \rangle = C \frac{1}{|x_1 - x_2|^{2\Delta}}$$

$$\mathcal{O}_i(x_1)\mathcal{O}_j(x_2) = \sum_k \frac{f_{ijk}}{|x_1 - x_2|^{\Delta_i + \Delta_j - \Delta_k}} \mathcal{O}_k(0)$$

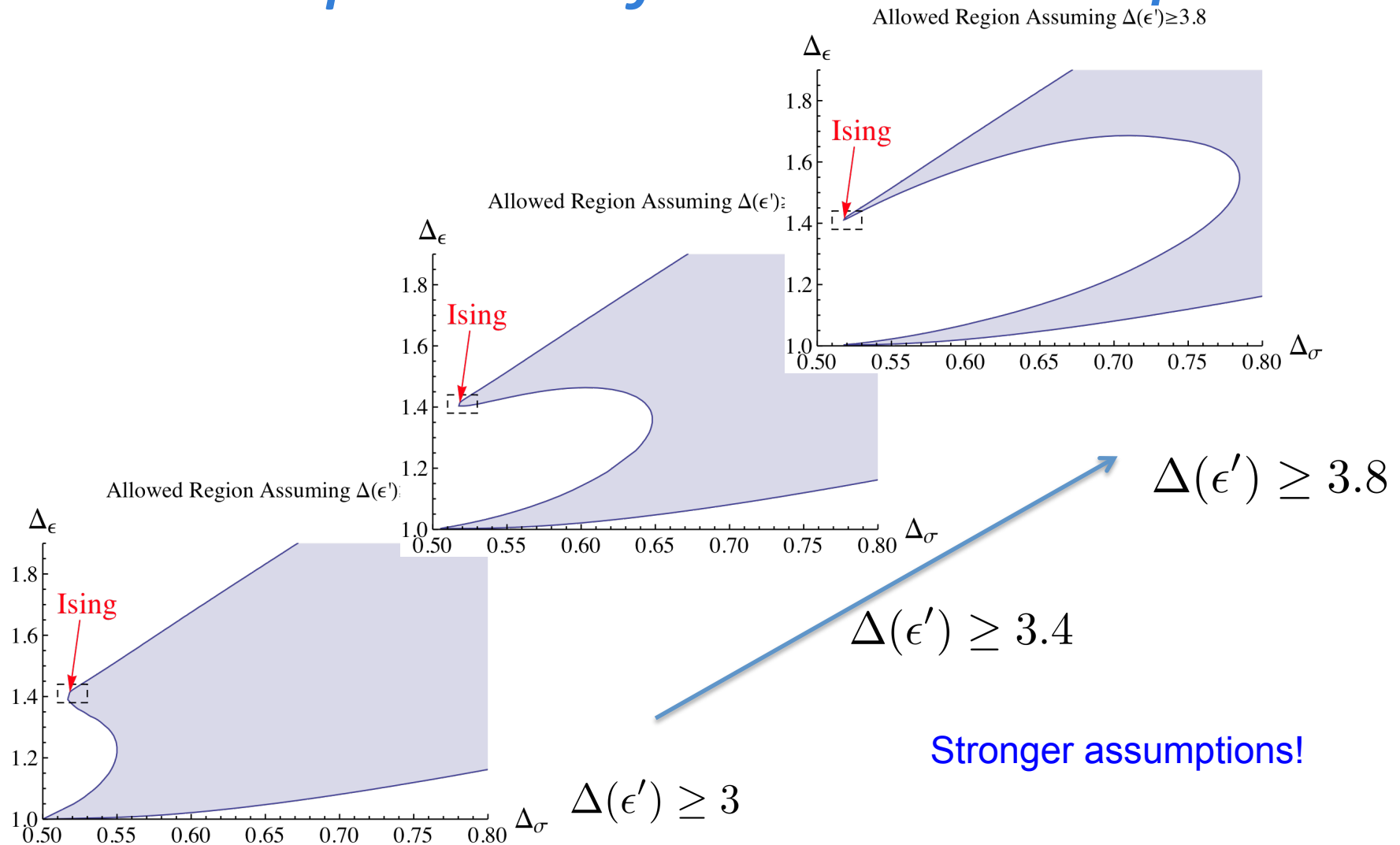
OPE & factorization completely fixed the theory\*  
(i.e. Data: spectral + couplings to conformal blocks)



complete sum over  
the conformal blocks  
“partial waves”.  
Only “tree” diagrams!

\* “Solving the 3D Ising Model with the Conformal Bootstrap” (El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin and Vichi) arXiv:1203.6064v1v [hep-th] (2012)

# Inequalities from Bootstrap\*



•“Solving the 3D Ising Model with the Conformal Bootstrap” (El-Showk, Paulos, Poland, Rychkov, Simmons-Duffin and Vichi) arXiv:1203.6064v1v [hep-th] (2012)

# Radial Quantization

Evolution:  $H = P_0$  in  $t \implies D$  in  $\tau = \log(r)$

$$ds^2 = dx^\mu dx_\mu = e^{2\tau} [d\tau^2 + d\Omega^2]$$

Can drop  
Weyl factor!

$$\mathbb{R}^d \rightarrow \mathbb{R} \times \mathbb{S}^{d-1}$$

"time"  $\tau = \log(r)$ , "mass"  $\Delta = d/2 - 1 + \eta$

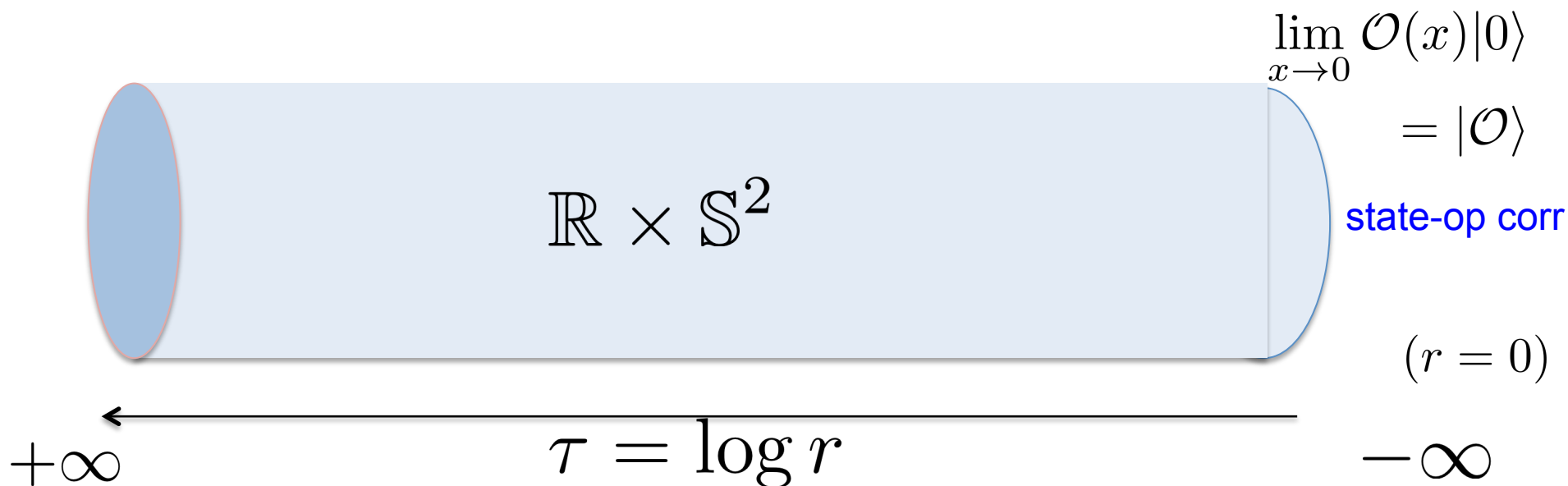
Exponentially Large Lattice:

$$a < \Delta r < L \quad \text{vs} \quad a < \Delta \log(r) < L$$



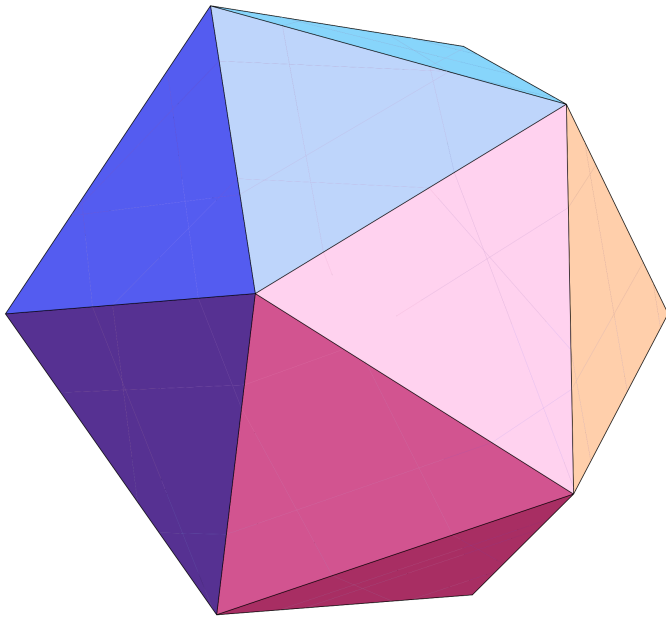
## 3-d Ising at Wilson-Fisher FP

$$Z_{Ising} = \sum_{\sigma(x,t)=\pm 1} e^{\beta \sum_{t, \langle x,y \rangle} \sigma(t,x)\sigma(t,y) + \beta \sum_{t,x} \sigma(t+1,x)\sigma(t,x)}$$

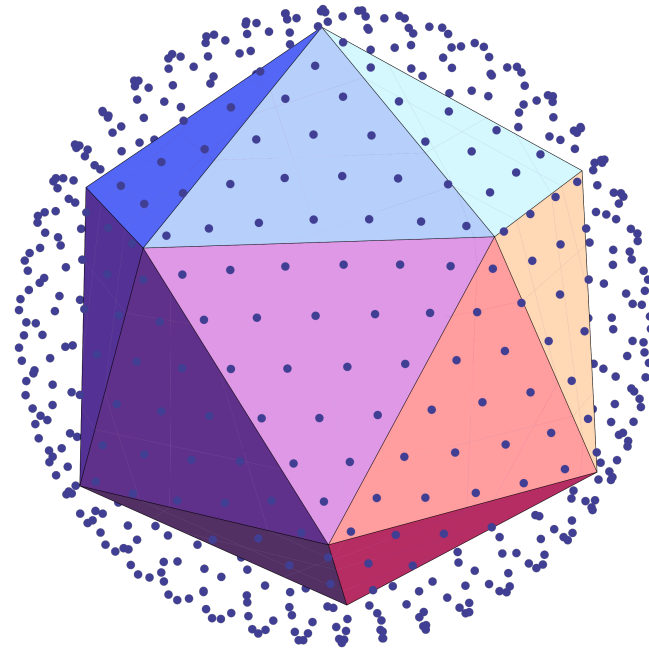


# *Order $s$ Refined Triangulated Icosahedron*

$s = 1$

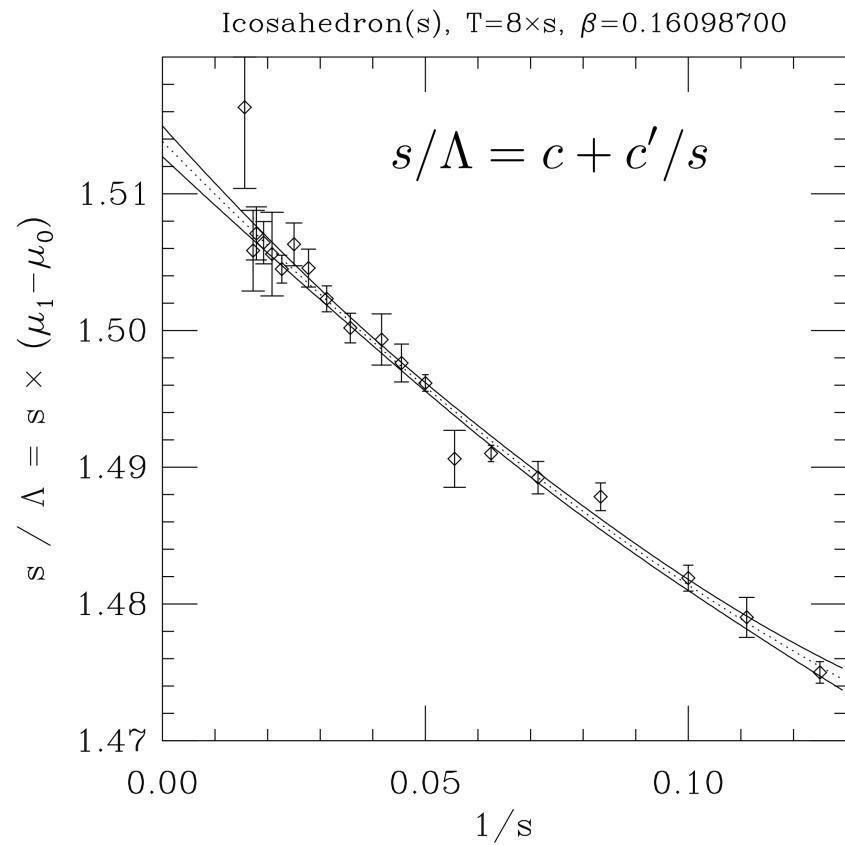
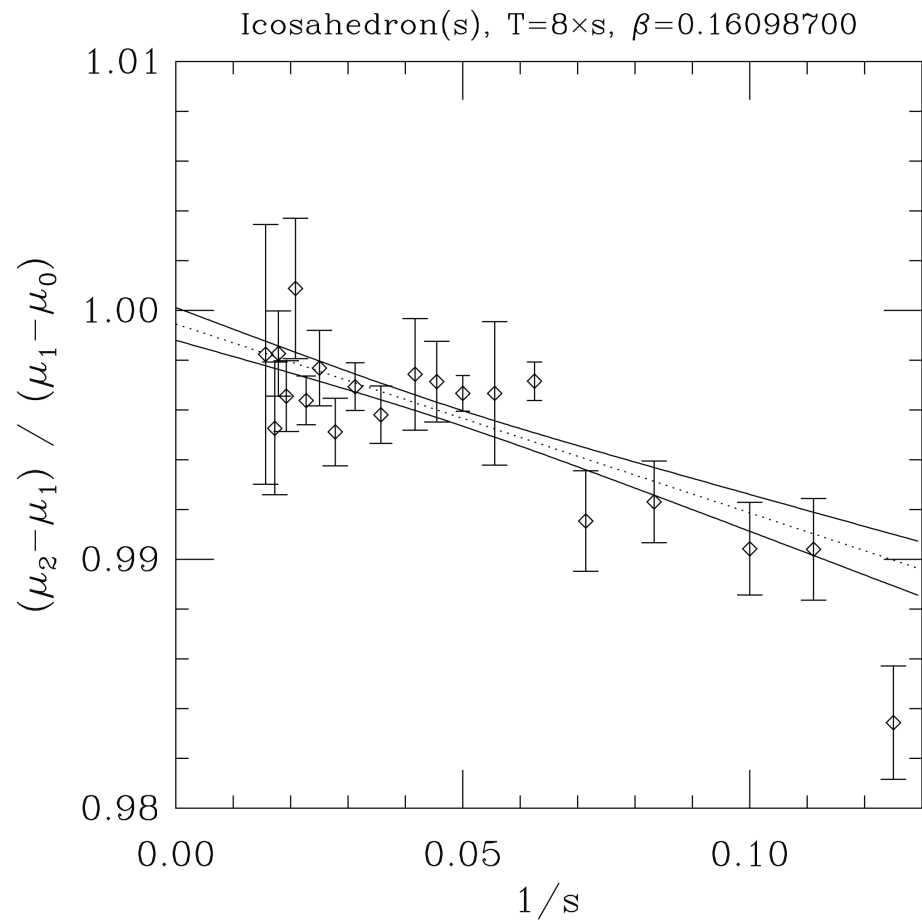


$s = 8$



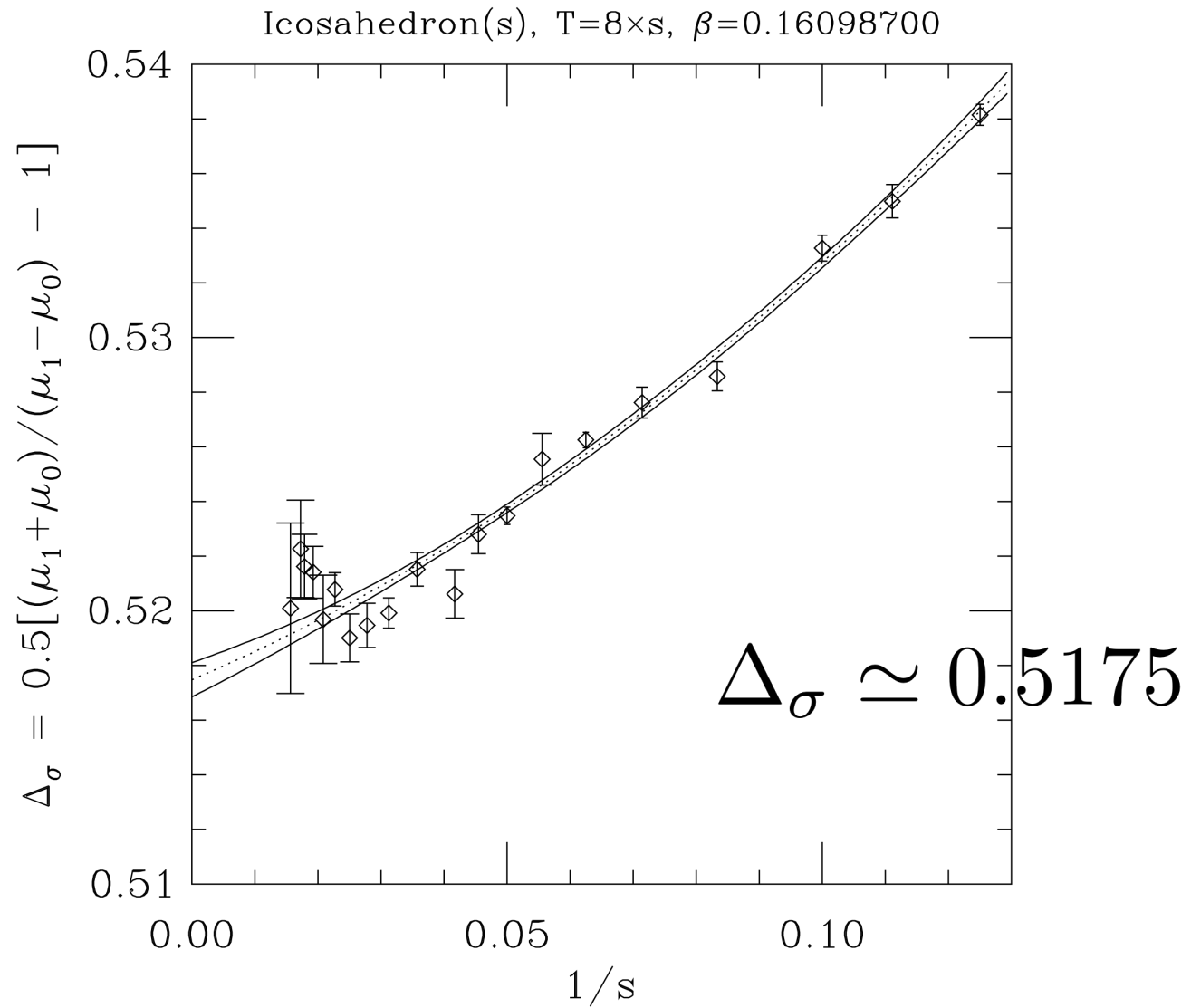
$l = 0$  (A),  $1$  (T1),  $2$  (H) are irreducible 120  
Icosahedral subgroup of  $O(3)$

# Check Descendant Relation & rescale “log(r)”



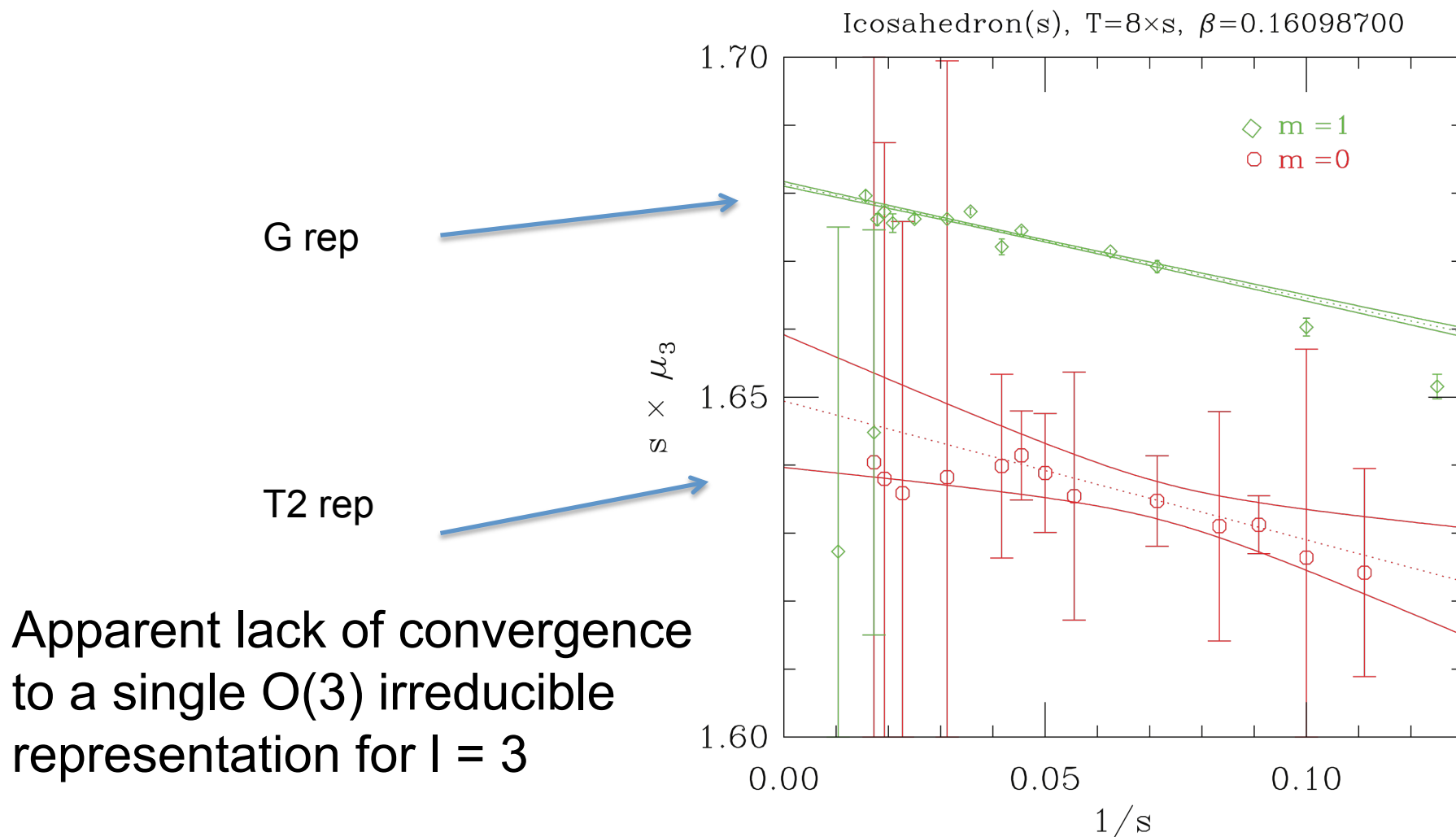
$$c = 1.5105(7)$$

# Current Fit:



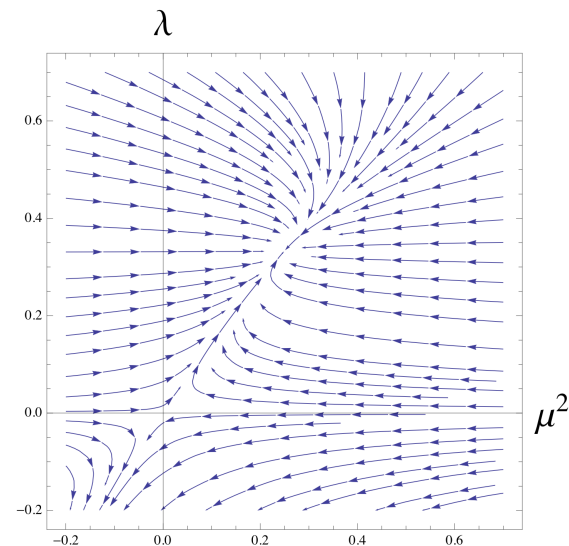
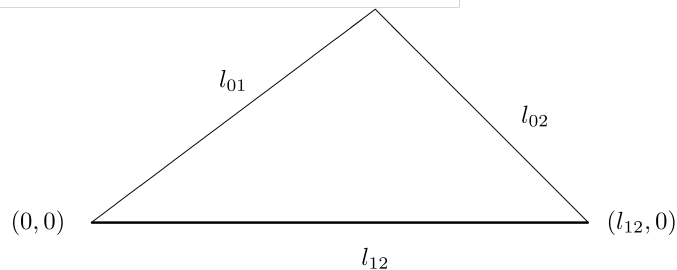
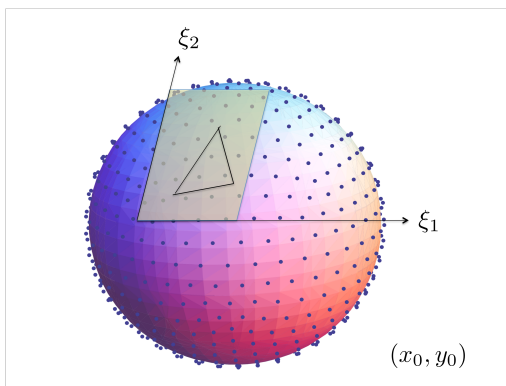
# Wrong Theory?

## Failure to recover $O(4,1)$ at $l = 3$ ?



# Finite Element Method/Regge Geometry

$$S = \int d^D x \sqrt{-g} \left[ \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi + \lambda \left( \phi^2 - \frac{\mu^2}{2\lambda} \right)^2 \right].$$



$$\int_{A_{012}} dx dy \partial_\mu \phi(x, y) \partial_\mu \phi(x, y) = \frac{1}{2A_{012}} [(l_{01}^2 + l_{20}^2 - l_{12}^2)(\phi_1 - \phi_2)^2 + \text{cyclic}]$$

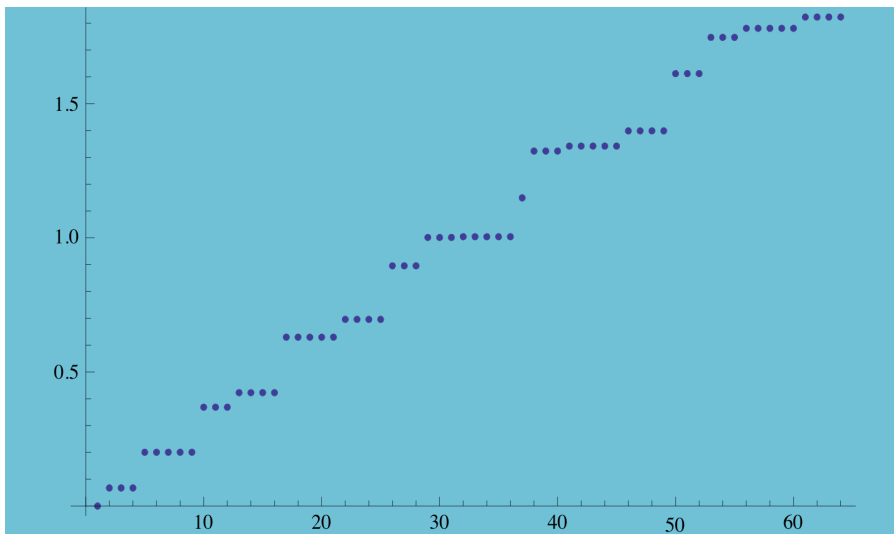
See also: Christ, Friedberg, Lee on "Random Lattice" NP (1982)

# *FEM fixes the huge Spectral defects*

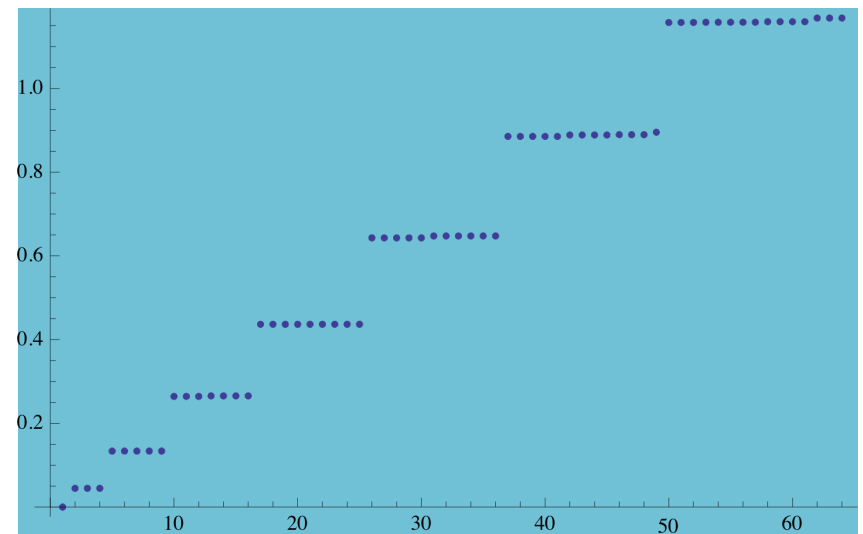
For  $s = 8$  first  $(l+1)*(l+1) = 64$  ev

BEFORE (K = 1)

AFTER (FEM K's)

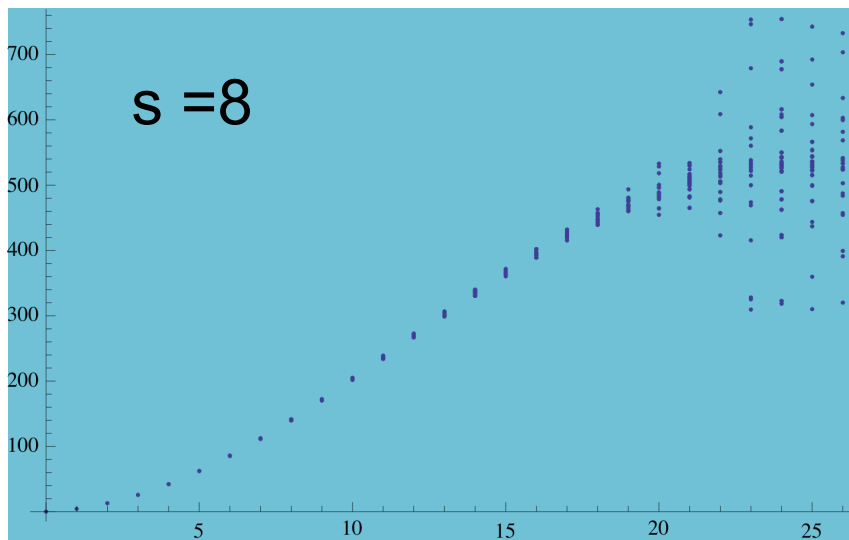
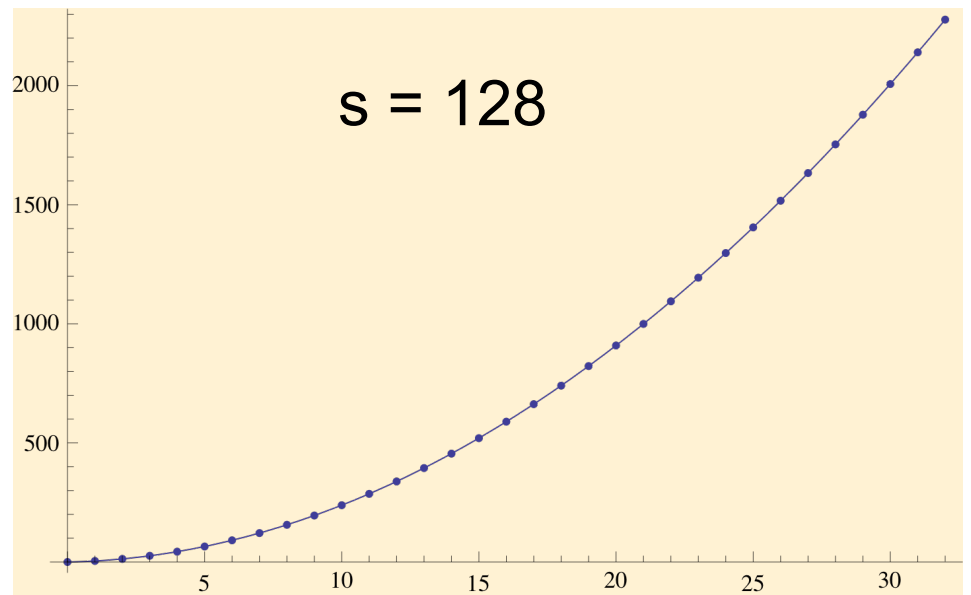
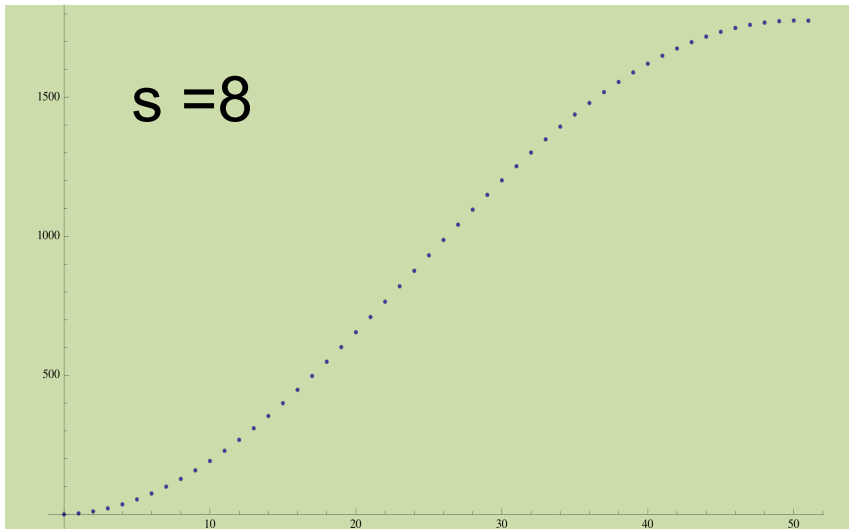


$l, m$



$l, m$

# *Spectrum of FE Laplacian on a sphere*



Fit

$$l + 1.00012 l^2$$

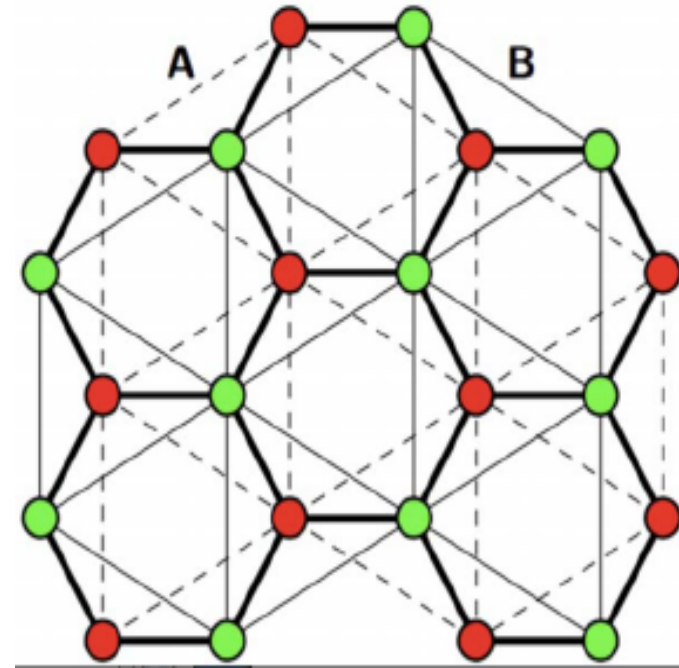
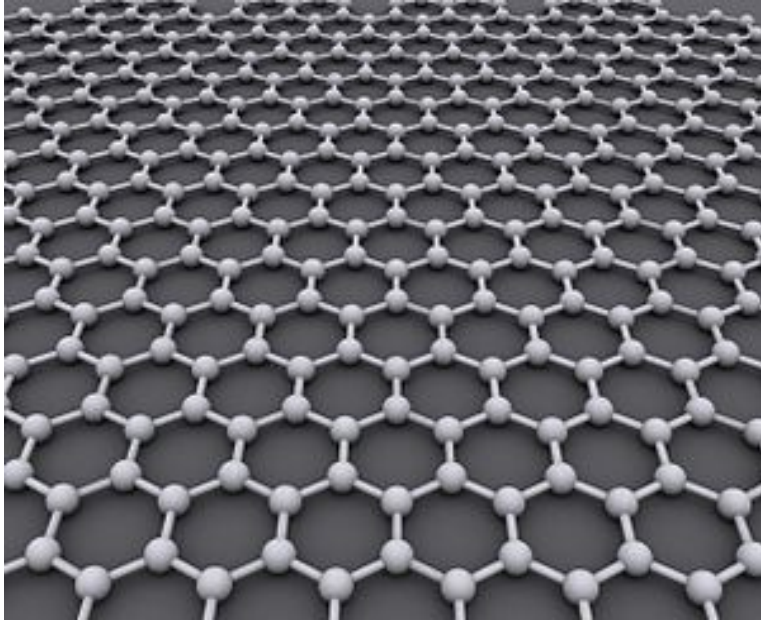
$$- 13.428110^{-6} l^3 - 5.5724410^{-6} l^4$$







# Graphene



2d Carbon hexagonal lattice: A/B Bravais sub lattices  $\sqrt{3}a$ ,  $a = 1.42\text{\AA}$

Effective field 2+1 relativistic theory: 4 copies of 2 component Dirac fields

Strong electrostatic fields:  $\longrightarrow e_{eff}^2 = \frac{e^2}{\hbar v} \simeq 300 \times \frac{e^2}{\hbar c}$

Phonons act like gauge fields.

# Tight Binding Hamiltonian

$$H = H_2 + H_C = a_{xs}^\dagger K_{xy} a_{ys} + q_x V_{xy} q_y$$

$$q_x = a_{x\uparrow}^\dagger a_{x\uparrow} + a_{x\downarrow}^\dagger a_{x\downarrow} - 1 \equiv a_x^\dagger a_x - b_x^\dagger b_x$$

*Exact internal symmetry:  $U(1) \times SU(2)$  Fermion number and Spin symmetry*

**Generators:**

$$Q = \sum_x q_x \rightarrow Q = [a_x^\dagger a_x - b_x^\dagger b_x]$$

$$J_{\pm} = a_{x,s}^\dagger \sigma_{\pm}^{ss'} a_{x,s'} \rightarrow J_+ = J_-^\dagger = (-1)^{B_x} a_x^\dagger b_x^\dagger$$

$$J_3 = a_{x,s}^\dagger \sigma_3^{ss'} a_{x,s'} / 2 \rightarrow J_3 = [a_x^\dagger a_x + b_x^\dagger b_x] / 2 - 1$$

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$s = \uparrow, \downarrow$  “a = up spin elect/b =down spin holes”

# Lagrangian for Tight Binding Hamiltonian

$$\mathcal{L} = \psi_x^\dagger (\partial_t + ie\phi\sigma_3 + e^2 V_{xx}) \psi_x - \kappa \sum_{\langle x,y \rangle} \psi_x^\dagger \psi_y + \frac{1}{4} \phi_x(t) V_{xy}^{-1} \phi_y(t)$$

Normal Ordering  
Term

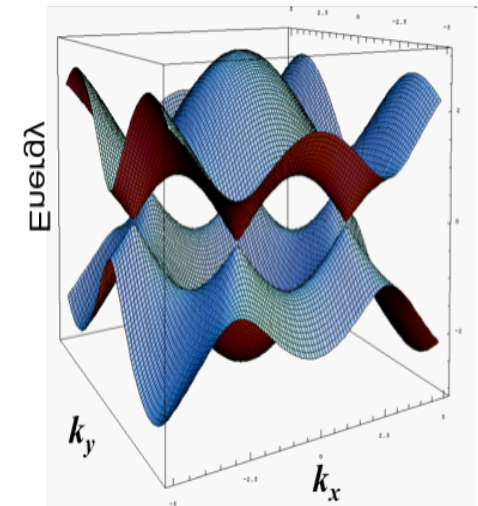
Invert and convolute by  
FFTs

Discretize time axis: Each spin has kinetic term (e=0)

$$D(E, k) = \begin{bmatrix} \overset{\text{A}}{(e^{iE\delta} - 1) + m\delta} & \overset{\text{B}}{\kappa c_k \delta} \\ \kappa c_k^* \delta & (e^{iE\delta} - 1) - m\delta \end{bmatrix} \begin{matrix} \text{A} \\ \text{B} \end{matrix}$$

where

$$\begin{aligned} \omega_k &= \pm \sqrt{c_k^* c_k} \\ &= \pm \sqrt{3 + 2 \cos(\sqrt{3}k_x a) + 4 \cos(\sqrt{3}k_y a/2) \cos(3k_x a/2)} \end{aligned}$$



# Recall Staggered Hypercubic Lattice

Unitary transform of Naive to Staggered Fermion

$$\sum_{\pm\mu} \bar{\psi}(x) \gamma_{\mu} \psi(x + \mu) = \sum a_{\mu}^{\dagger} \eta_{\mu}(x) a_{x+\mu}$$

$$\text{where } \gamma_{-\mu} \equiv -\gamma_{\mu}$$

Defining the Unitary matrix

$$S(x_1, x_2, x_3, x_4) = \gamma_1^{x_1} \gamma_2^{x_2} \gamma_3^{x_3} \gamma_4^{x_4}$$

$$\psi_y = S(y) a_y \quad , \quad \psi_x^{\dagger} = a_x^{\dagger} S^{\dagger}(x)$$

$$\eta_{\mu}(x) = S^{\dagger}(x) \gamma_{\mu} S(x + \mu) = (-1)^{x_1 + \dots + x_{\mu-1}}$$

# Graphene Naive Dirac Equation

$$\hat{e}_1 + \hat{e}_2 + \hat{e}_3 = 0 \qquad \hat{e}_i \cdot \hat{e}_j = \frac{3\delta_{ij} - 1}{2}$$

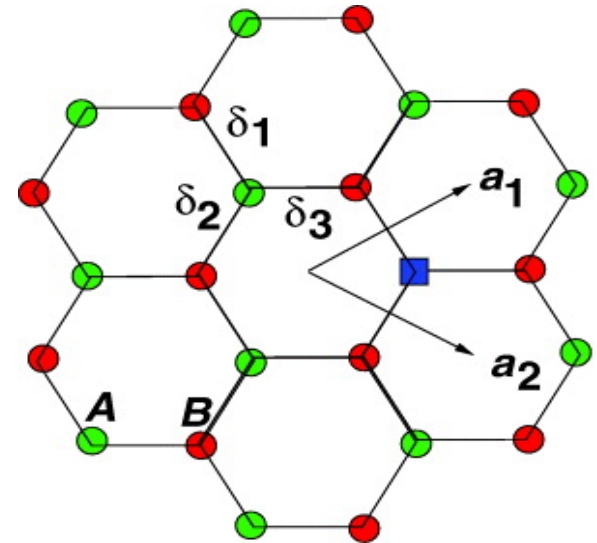
Defining:

$$\sigma^{(\mu)} \equiv \hat{e}_\mu \cdot \vec{\sigma} \implies \sigma^{(1)} + \sigma^{(2)} + \sigma^{(3)} = 0$$

## Graphene Dirac Equation

$$\sum_{\mu=1,2,3} \psi_x^\dagger \hat{e}_\mu \cdot \vec{\sigma} \psi_{x+\hat{e}_\mu} + h.c.$$

With  $x$  restricted to A (green) sublattice  
 h.c. give B (red) sublattice contribution  
**NOTE: Hermiticity require all backward link B to A also have + sign!**



# Graphene “Staggered” Transformation

Path ordered closed Path give +1. For example one Hexagon is

$$\sigma^{(1)}\sigma^{(2)}\sigma^{(3)}\sigma^{(1)}\sigma^{(2)}\sigma^{(3)} = -\sigma^{(1)}\sigma^{(3)} - \sigma^{(2)}\sigma^{(3)} = 1$$

This implies path independence for the path order spin product

$$S(x) = \sigma^{(\mu_k)} \dots \sigma^{(\mu_3)}\sigma^{(\mu_2)}\sigma^{(\mu_1)}$$

so we introduce local operators  $\psi_y = S(y)a_y$  ,  $\psi_x^\dagger = a_x^\dagger S^\dagger(x)$

to show equivalence to the traditional 2 spin tight binding action

$$\sum_{\mu} \psi_x^\dagger \hat{e}_{\mu} \cdot \vec{\sigma} \psi_{x+\hat{e}_{\mu}} + h.c. = \sum_{\mu} [a_{\uparrow,x}^\dagger a_{\uparrow,x+\hat{e}_{\mu}} + a_{\downarrow,x}^\dagger a_{\downarrow,x+\hat{e}_{\mu}}] + h.c.$$

# Improved TIME DISCRETIZATION ?

Continuum D.R.  $E = \pm i \sqrt{\omega_k^2 + m^2}$  with  $\omega_k^2 = c_k^* c_k$

Forward  
Differences form

$$2 \sin \delta E / 2 = \pm i \delta e^{-i \delta E / 2} \sqrt{\omega_k^2 + m^2}$$

$O(\delta)$  Errors

A/B Time  
reversal form

$$2 \sin E \delta / 2 = \pm i \delta \sqrt{\frac{m^2 + \omega_k^2}{1 - m}}$$

$O(\delta^2)$  Errors

Staggered  
Time form

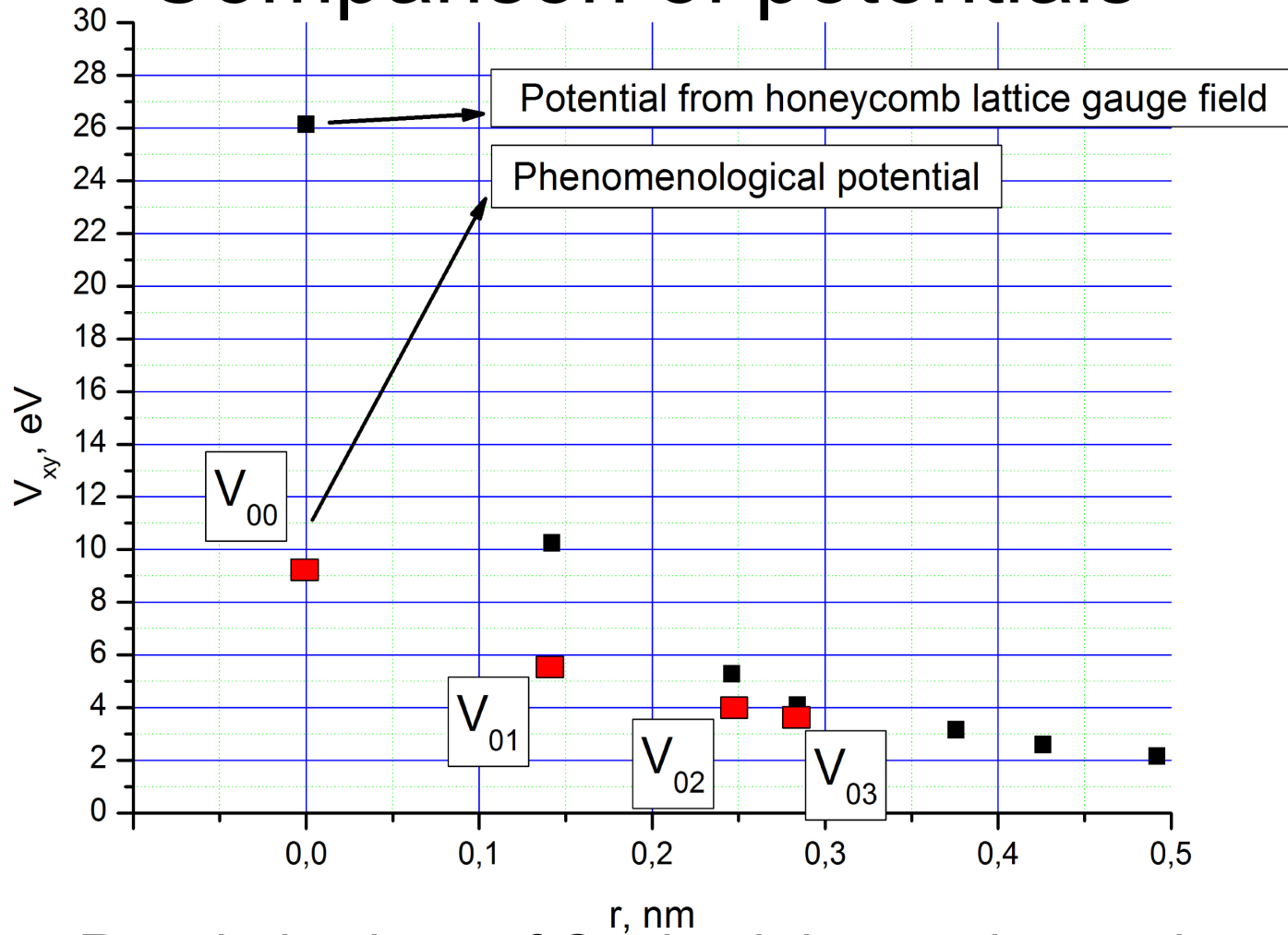
$$\sin \delta E = \pm \delta i \sqrt{\omega_k^2 + m^2}$$

$O(\delta^2)$  Errors

Comparison of 3 options need testing on realistic lattices  
with Coulomb term

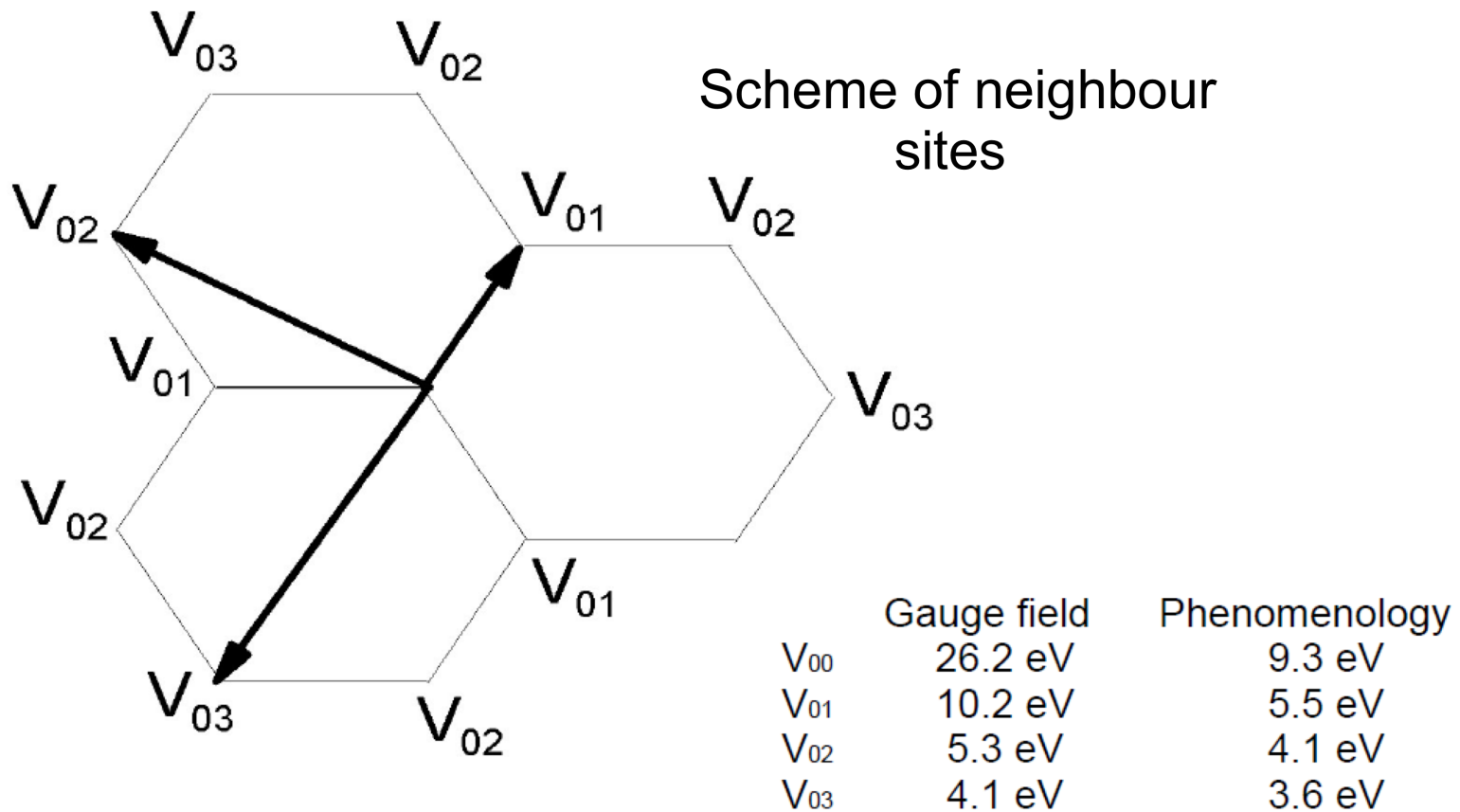


# Comparison of potentials



«Regularization» of Coulomb interaction at short distances

# Comparison of potentials



see M. V. Ulybyshev, P. V. Buividovich, M. I. Katsnelson, ArXiv: 1304.3660, to appear in PRL

# Conclusions

Regularization of Coulomb potential at short distances strongly affects the conductor-insulator phase transition in graphene.

Though the suspended graphene is a conductor, the phase transition still exists at unphysical dielectric permittivity  $\epsilon \sim 0.7$ .

M. V. Ulybyshev, P. V. Buividovich, M. I. Katsnelson, ArXiv: 1304.3660, to appear in PRL

<http://www.yale.edu/QCDNA/index14.html> June 19-21, 2014